# **PHI IN THE SKY 2004**

# Linearized Bekenstein Varying Alpha Models

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# **Bekenstein Type Model**

Bekenstein Type Model

Linearized case

• Equations of motion

Predictions

Predictions

Predictions

Conclusions

 $\bullet e = e\left(\phi(x^{\mu})\right)$ 

(1)

$$\mathcal{L} = \mathcal{L}_{\phi} + \mathcal{L}_{\phi F} + \mathcal{L}_{\mathsf{other}}$$

(2) 
$$\mathcal{L}_{\phi} = \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - V(\phi) ,$$

3) 
$$\mathcal{L}_{\phi F} = -\frac{1}{4} B_F(\phi) F_{\mu\nu} F^{\mu\nu} = -\frac{\alpha_0}{4\alpha} F_{\mu\nu} F^{\mu\nu} ,$$



## **Linearized case**

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### Linear coupling

Assuming small change in  $\phi$ :

$$B_F(\phi) = 1 + \zeta_F(\phi - \phi_0) + \frac{1}{2}\xi_F^2(\phi - \phi_0)^2 + \dots$$

$$-\frac{\Delta \alpha}{\alpha} = \zeta_F \left(\phi - \phi_0\right) + \frac{1}{2} \left(\xi_F - 2\zeta_F\right) \left(\phi - \phi_0\right)^2 + \dots$$



## **Linearized case**

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### Linear Potential

 $V(\phi), \alpha(\phi)$  always linear functions of  $\phi$  around today for limited period of time  $\Rightarrow$  few free parameters

(4) 
$$V(\phi) = V(\phi_0) + \frac{dV}{d\phi} (\phi - \phi_0), \quad \frac{dV}{d\phi} < 0$$

(5) 
$$\alpha = \alpha_0 + \frac{d\alpha}{d\phi} \left(\phi - \phi_0\right) \,,$$



# **Equations of motion**

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### • $\phi$ source of dark energy

$$H^{2} = H_{0}^{2} \left( \Omega_{m0} a^{-3} + \Omega_{r0} a^{-4} + \Omega_{\phi} \right) ,$$

(7) 
$$\frac{\ddot{a}}{a} = -H_0^2 \left[ \frac{\Omega_{m0}}{2} a^{-3} + \Omega_{r0} a^{-4} + \frac{\Omega_{\phi}}{2} (1+3w_{\phi}) \right] ,$$

where

(6)

(8) 
$$\Omega_{\phi} = \frac{8\pi(\dot{\phi}^2/2 + V(\phi))}{3H_0^2} ,$$

and

(9)

$$\omega_{\phi} = \frac{\dot{\phi}^2/2 - V(\phi)}{\dot{\phi}^2/2 + V(\phi)} \,.$$



# **Equations of motion**

Bekenstein Type Model

(10)

Linearized case

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$$\ddot{\phi} + 3H\dot{\phi} = -\frac{dV}{d\phi} + \frac{\alpha_0}{4\alpha^2}\frac{d\alpha}{d\phi}F_{\mu\nu}F^{\mu\nu}$$

If  $\frac{dV}{d\phi} = 0$  equivalence principle  $\Rightarrow |\zeta_F| < 5 \times 10^{-4} \Rightarrow \frac{\Delta \alpha}{\alpha} \simeq 10^{-10}$  (too small!)

Generalized Bekenstein models:  $\phi$  is driven by dark matter couplings [1],  $\phi$  is driven by its own potential [2]

 $\dot{\phi_0}, V_0, \frac{dV}{d\phi}, \frac{d\alpha}{d\phi}$  + cosmological parameters  $\Rightarrow$ 

### complete description of a particular linearized model

H. B. Sandvik, J.D. Barrow J. Magueijo, Phys. Rev. Lett **88**, 031302 (2002), hep-ph/0110377 E.J. Copeland, N.J. Nunes and M. Pospelov, Phys. Rev. **D69**, 023501 (2004), hep-ph/0307299

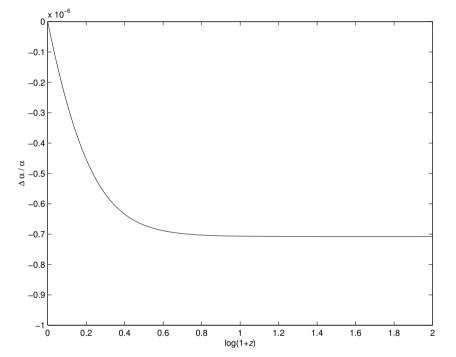




- Linearized case
- Equations of motion

#### Predictions

- Predictions
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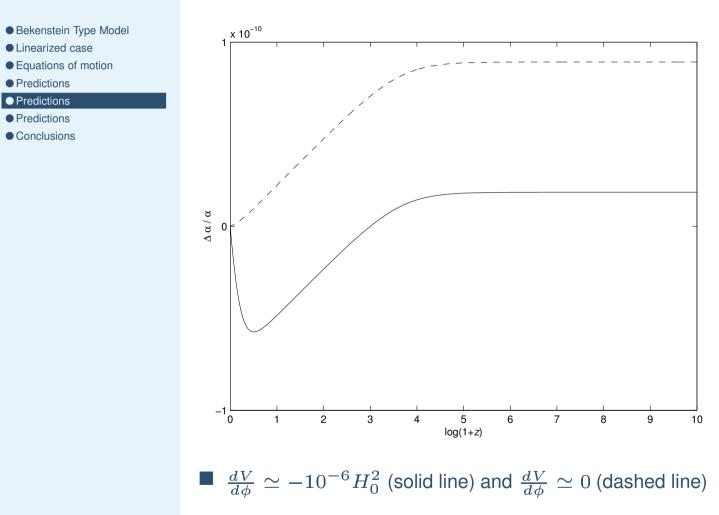


Unambiguously predicted by our model up to a normalization factor

(11) 
$$\zeta_{\alpha} \equiv -\frac{dV}{d\phi} \frac{\zeta_F}{H_0^2}$$

given  $\Omega_{\phi}^0 \simeq 0.7$  and  $\omega_{\phi}^0 \sim -1$ 







(1

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Constraints on  $\delta \phi \equiv \phi(z=0) - \phi(z=1)$ :

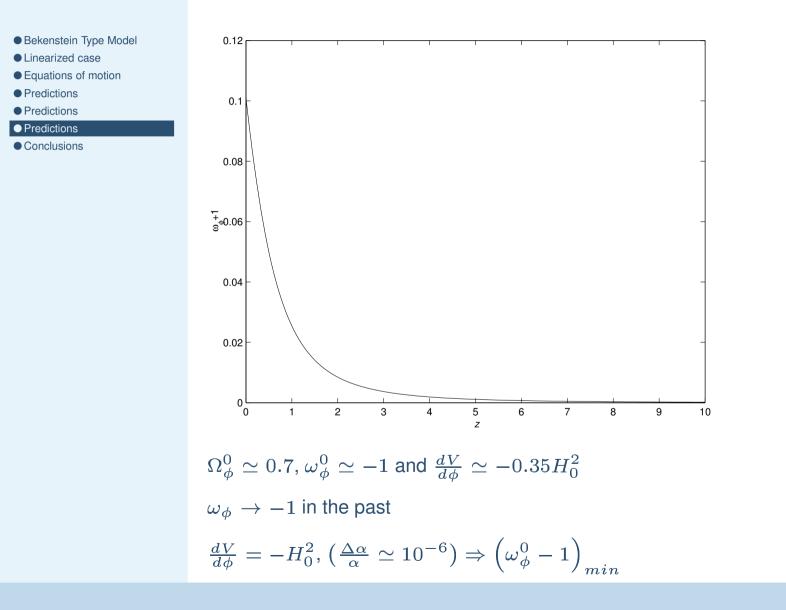
2) 
$$\delta V \le \rho_{c0} \Rightarrow \delta \phi \le 10^{-1}$$
.

Equivalence Principle  $\Rightarrow \delta \phi \geq 10^{-3}$  considering  $\frac{\delta \alpha}{\alpha} \sim 10^{-6}$ 

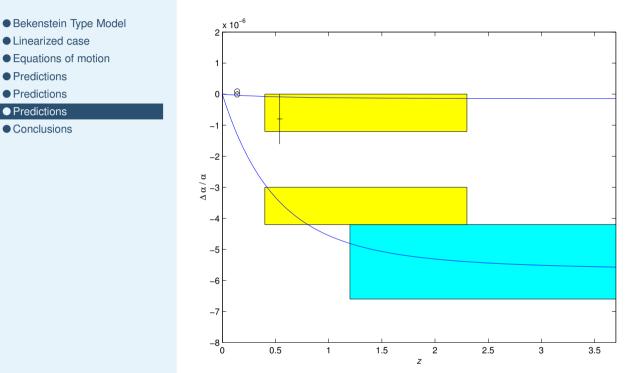
■  $10^{-3} \le \delta \phi \le 10^{-1}$ 

 $-H_0^2 \le \frac{dV}{d\phi} \le -10^{-2}H_0^2$ 









Oklo ([1], circles), Rhenium decay from meteorites ([2], vertical bar) and quasar data (Murphy *et al.* [3]) blue shaded box, Chand *et al.* ([4], yellow shaded boxes - up with terrestrial isotopic abundances and down with low-metalicity isotopic abundances)

Y. Fujii, Astrophysics. Space Sci 283, 559 (2003) gr-qc/0212017
K.A.Olive et al., Phys.Rev. D69,027701 (2004), astro-ph/0309252
M. T. Murphy, J. K. Webb, and V. V. Flambaum, Mon. Not. Roy Astron. Soc 345 609 (2003), astro-ph/0210299
H. Chand R. Srianand, P. Petijean, anf B. Aracil (2004), astro-ph/0210299



# Conclusions

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• We have studied the linearized class of Bekenstein-type varying  $\alpha$  models

Any realistic Bekenstein model should reduce to these models for a certain time interval around the present day and we are assuming a cosmological time

- Very specific predictions were obtained and compared with existing data:
  - No such linearized model is consistent with all existing observational results at different z

This could be explained by either unknown systematic errors in the observations or linearity breaking down for  $t \ll H^{-1}$