

Quantum Cosmology with a Chaplygin Gas

MB Lopez and PV Moniz

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- The theory $\leftarrow$ from Strings/Branes
- *Mise en Scene*: Chaplygin gas cosmology

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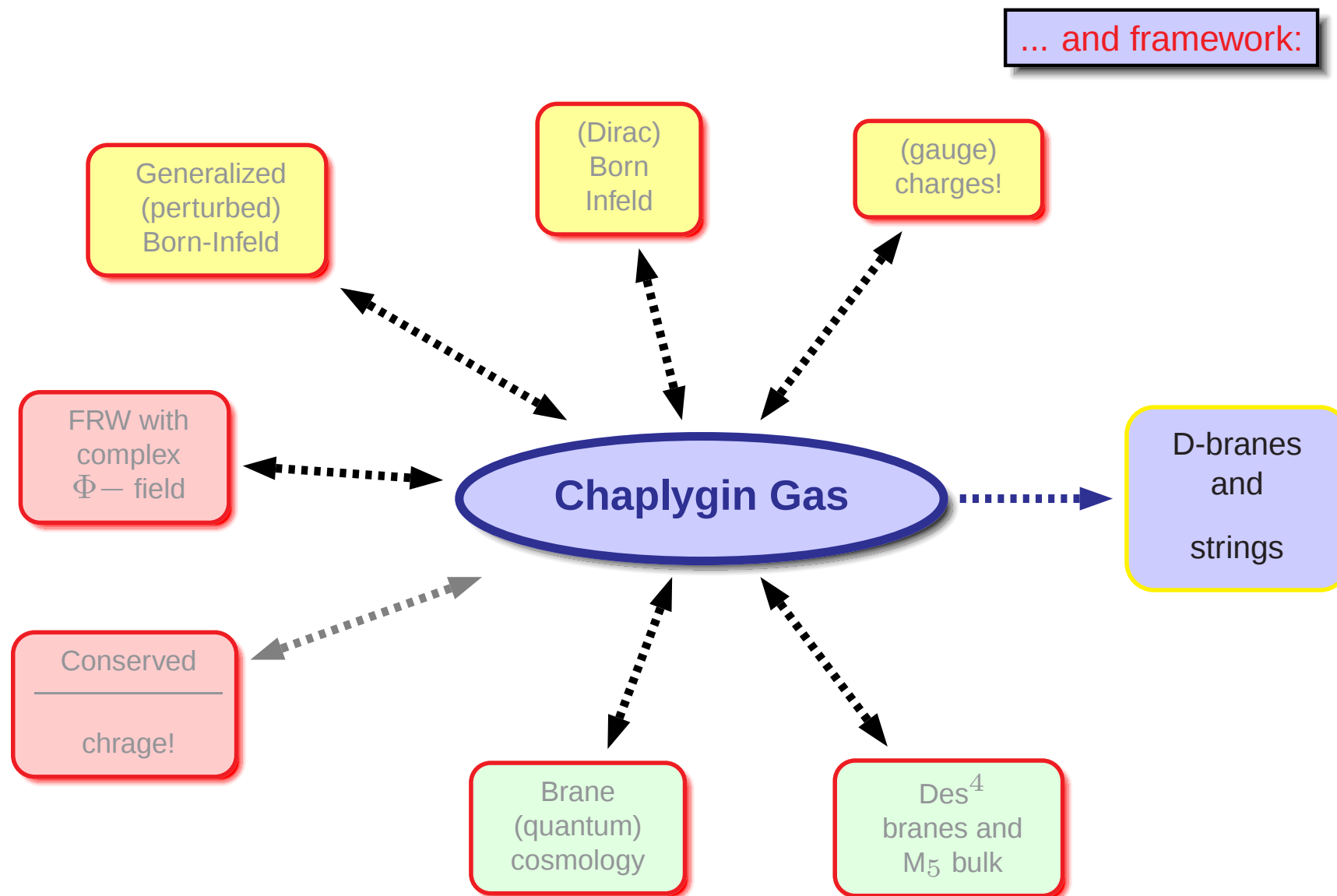
- The theory $\leftarrow$ from Strings/Branes
- *Mise en Scene*: Chaplygin gas cosmology
- A Physical Guide: solutions and quantization

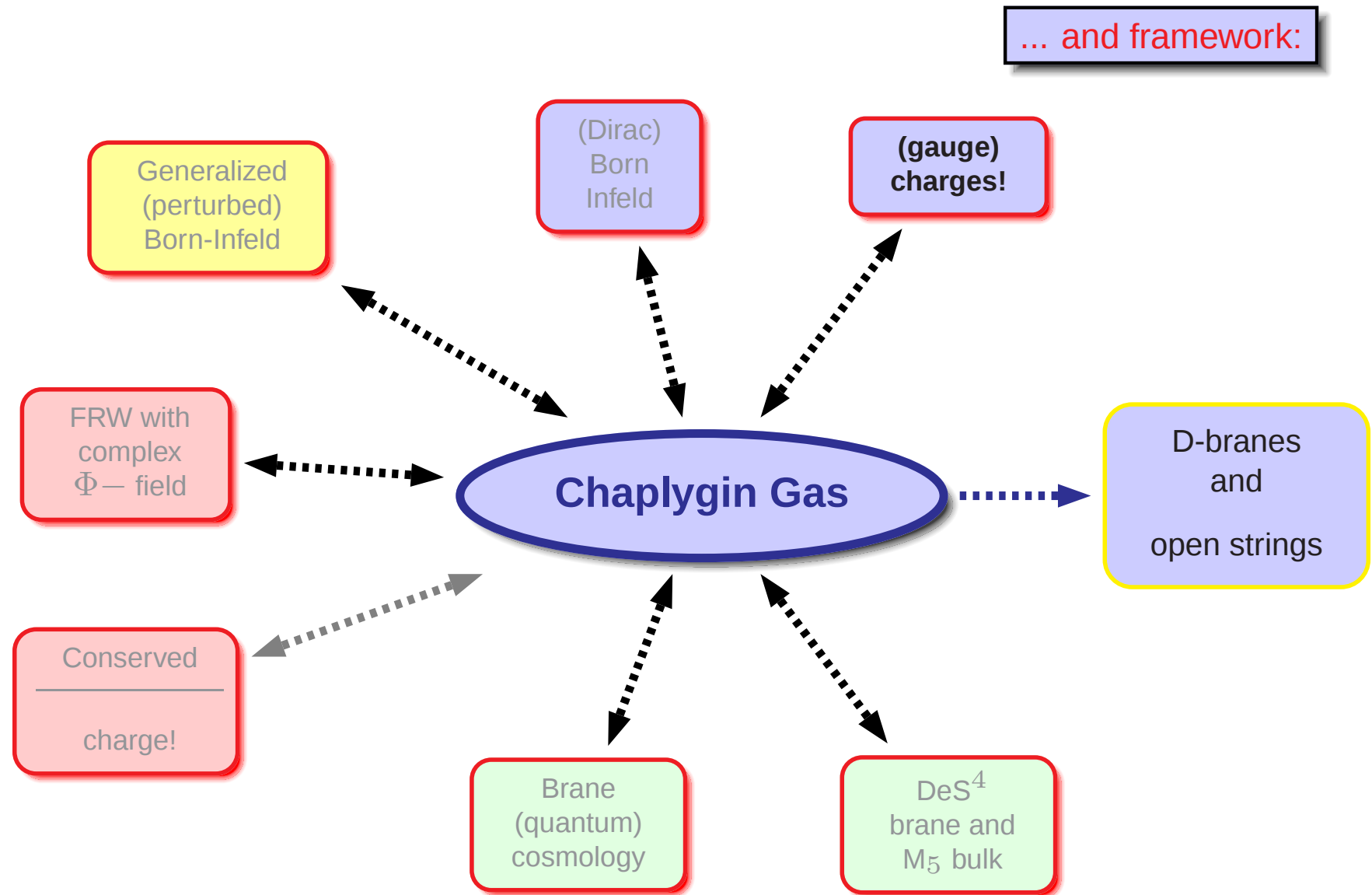
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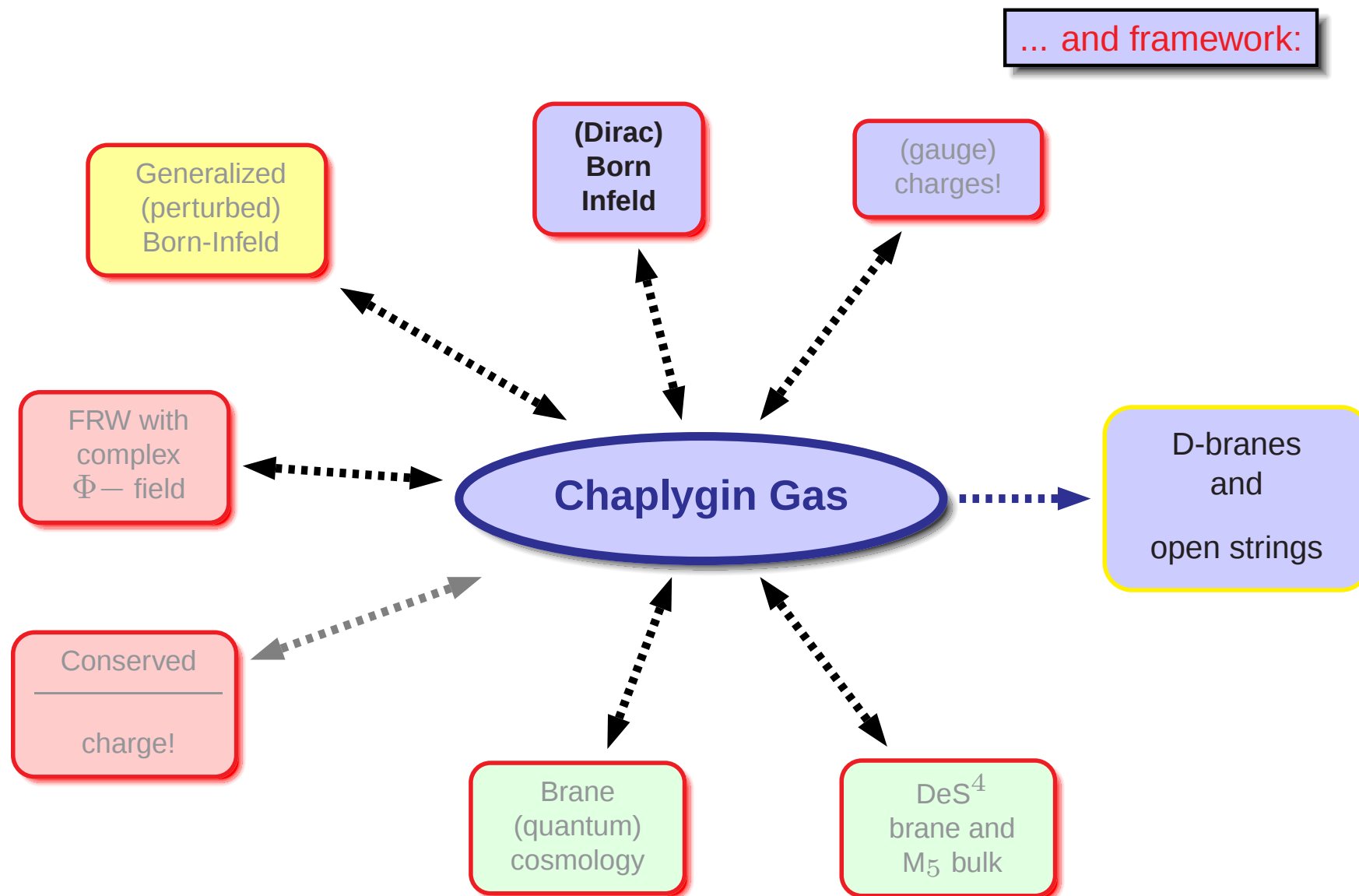
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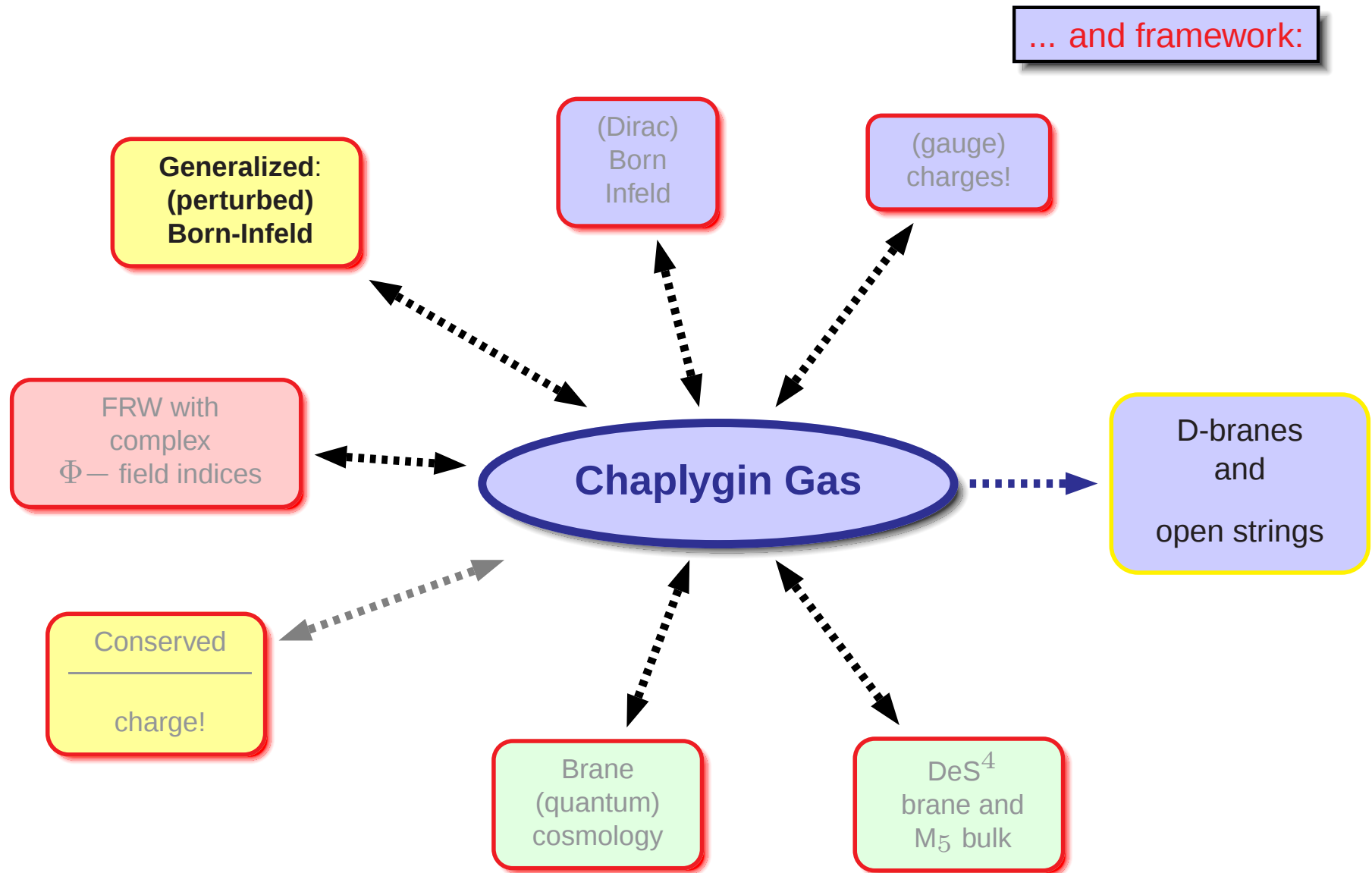
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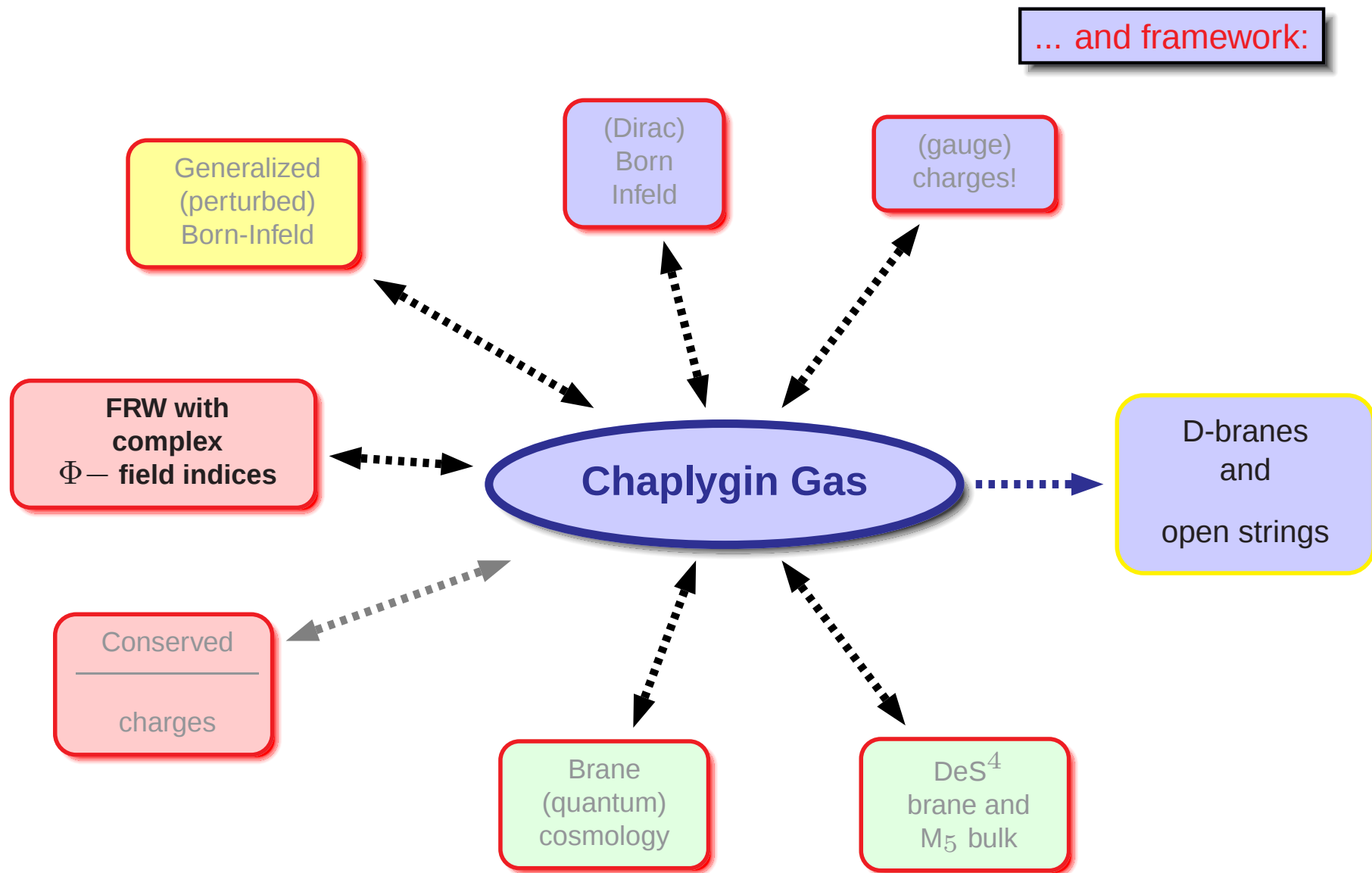
- The theory ← from Strings/Branes
- *Mise en Scene*: Chaplygin gas cosmology
- A Physical Guide: solutions and quantization
- *Finale* ... well, not yet ...

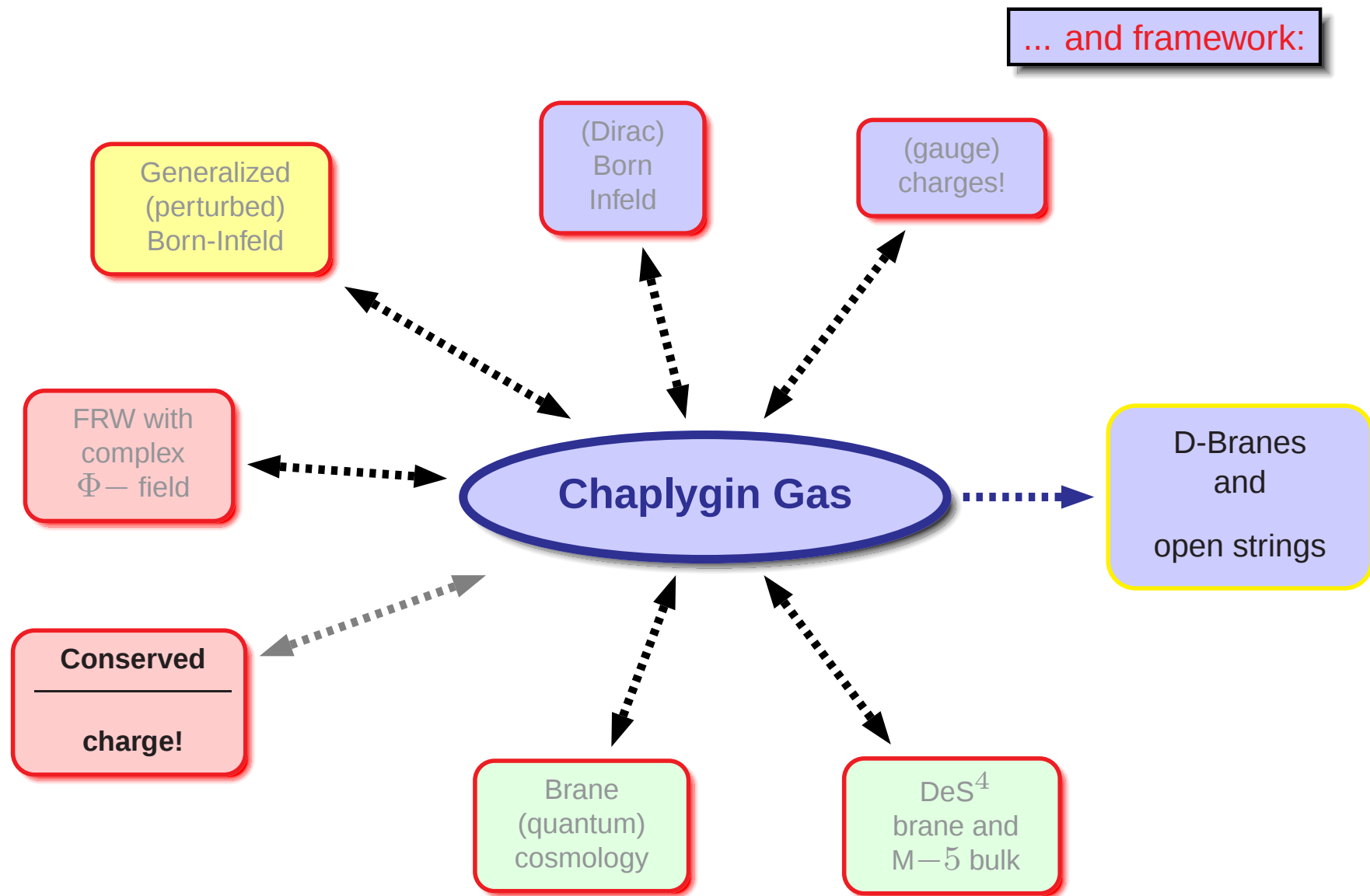


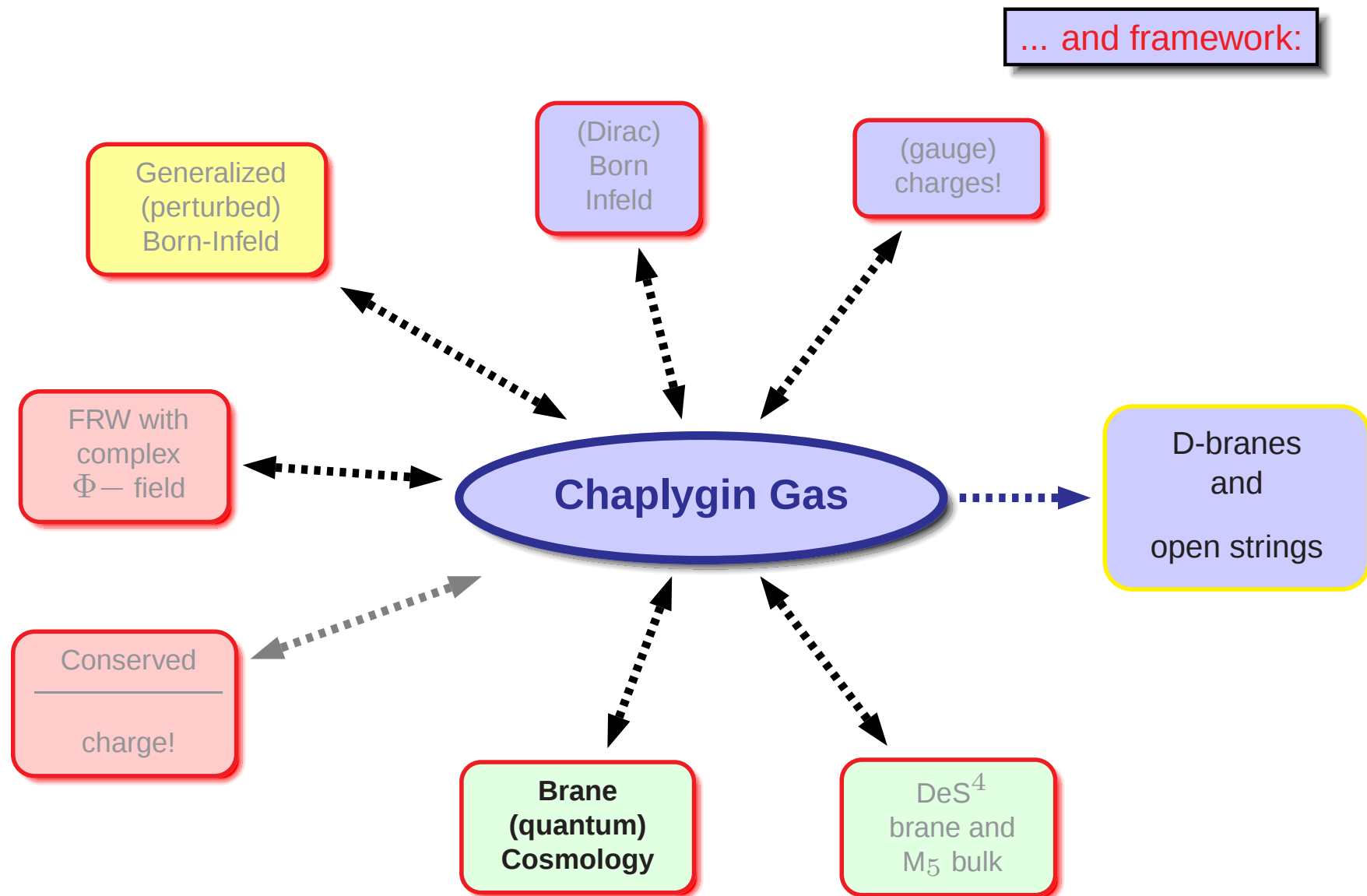


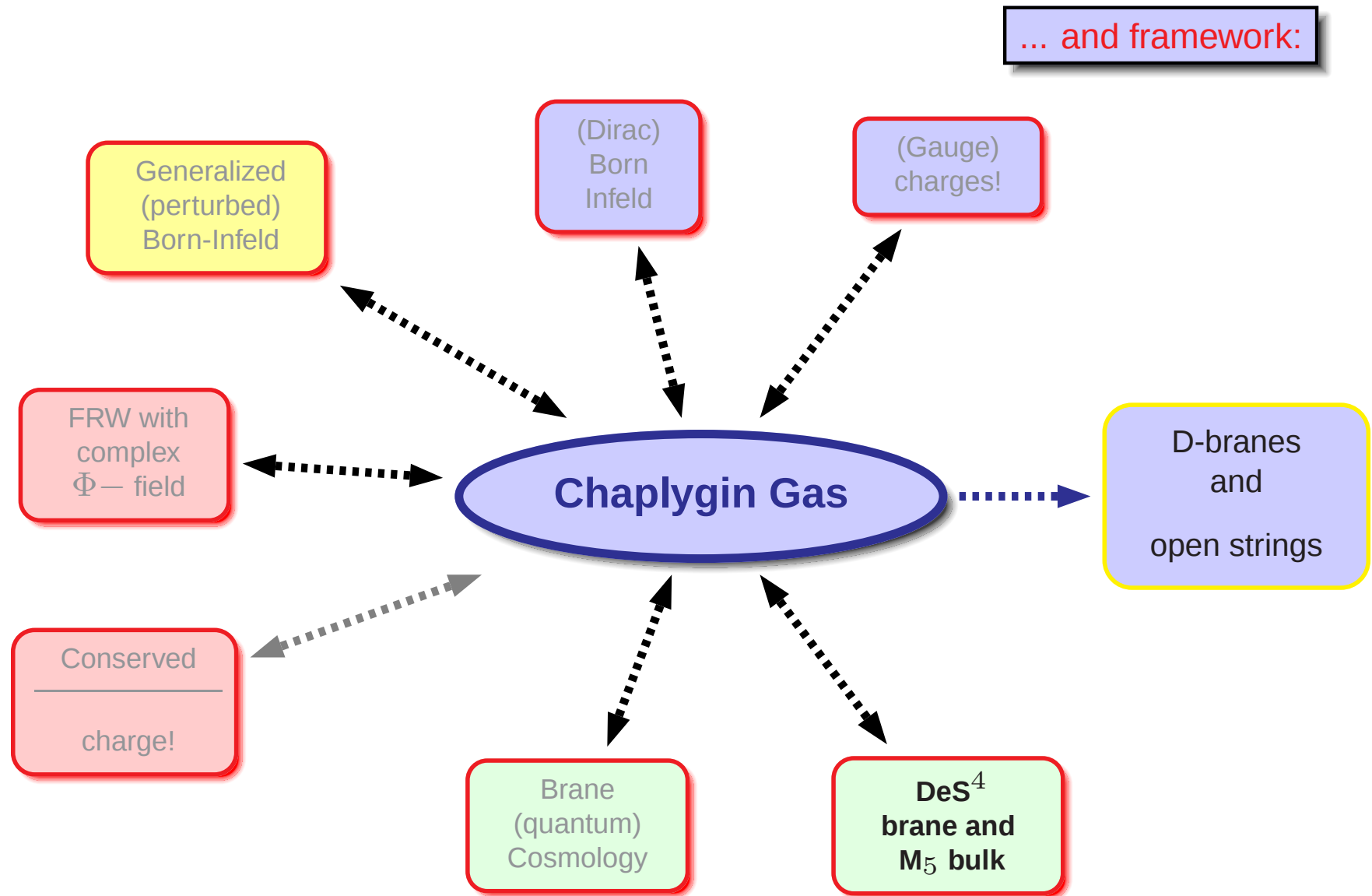




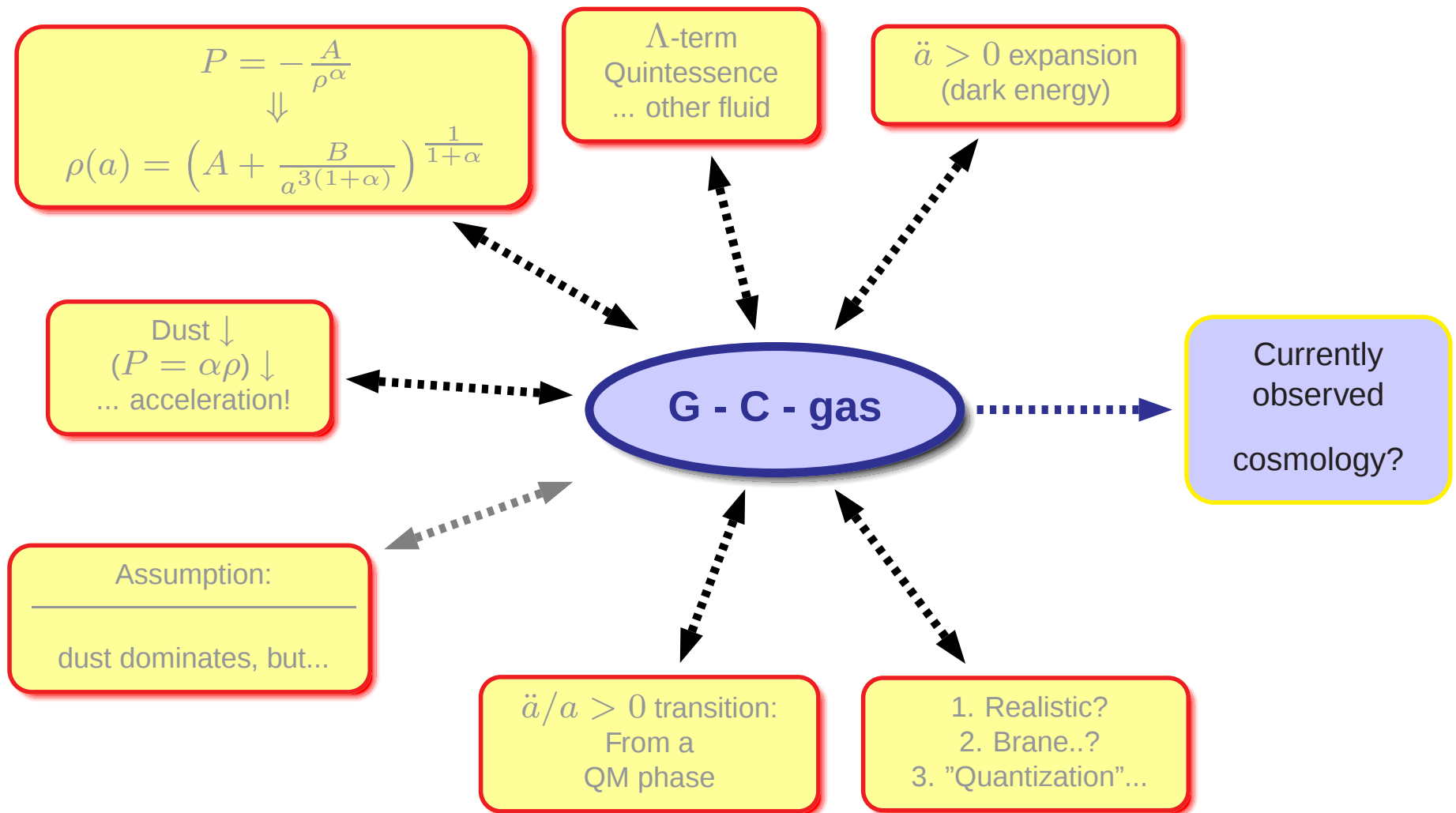




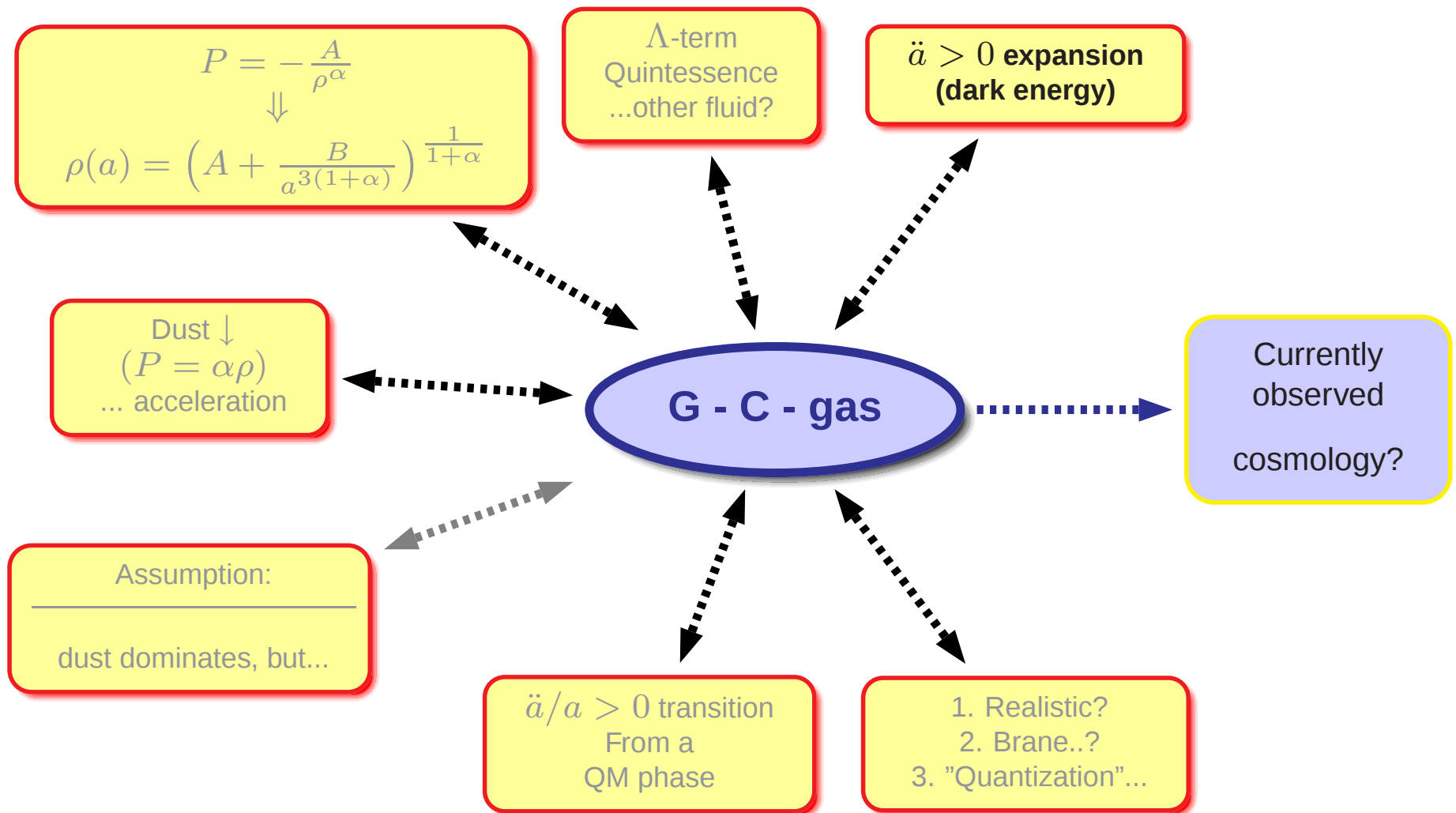




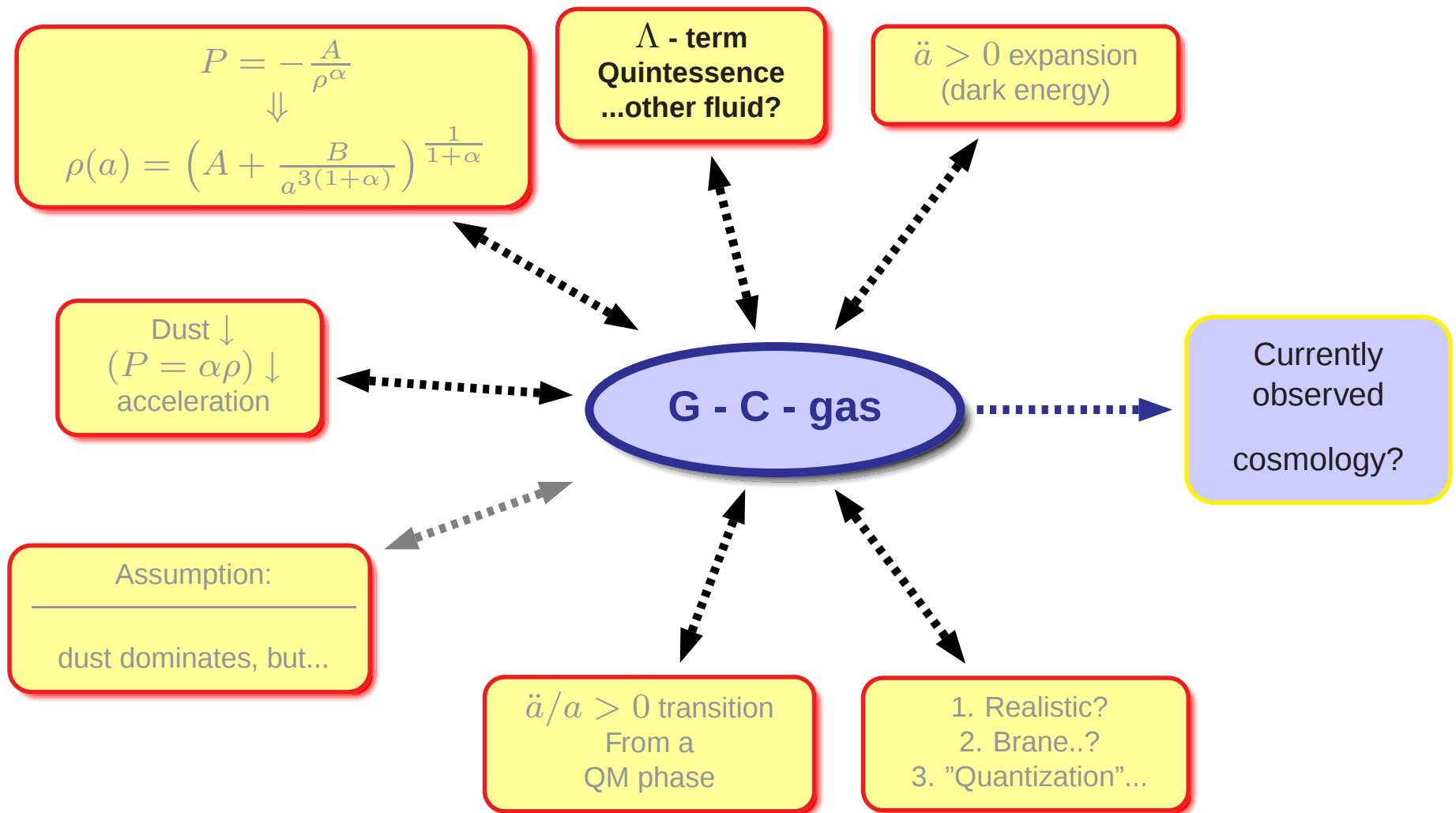
... and Mise en Scene



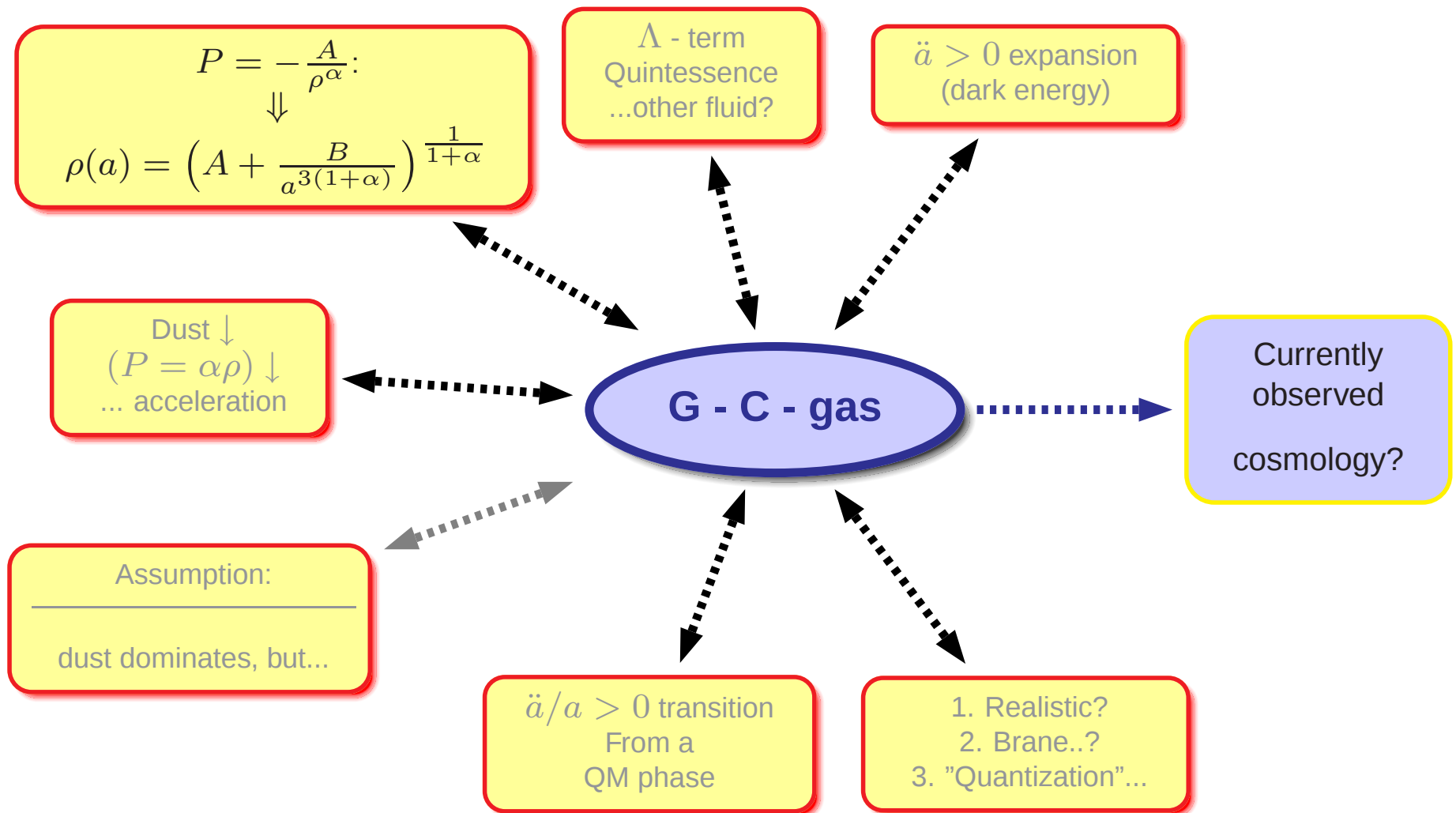
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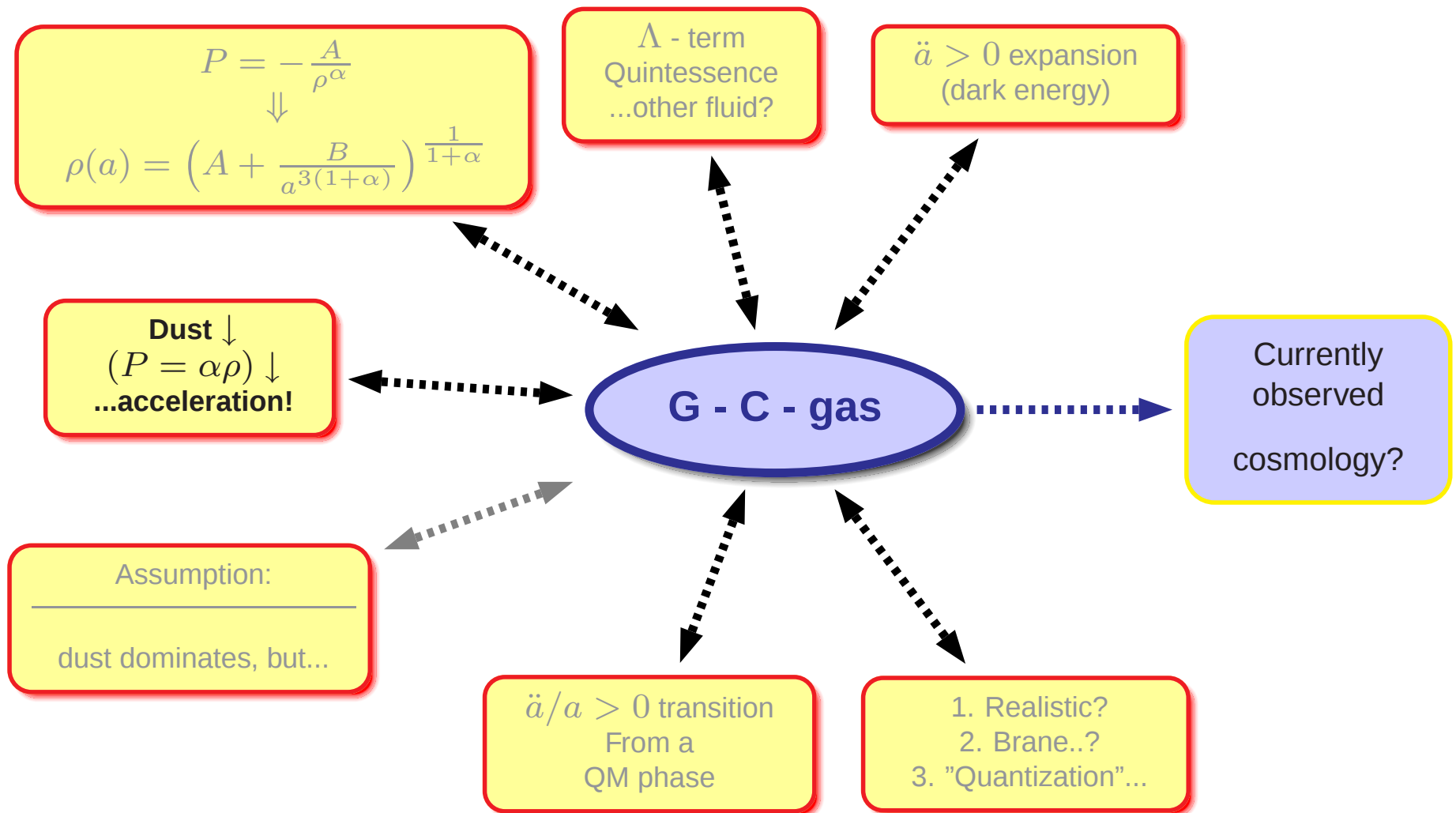
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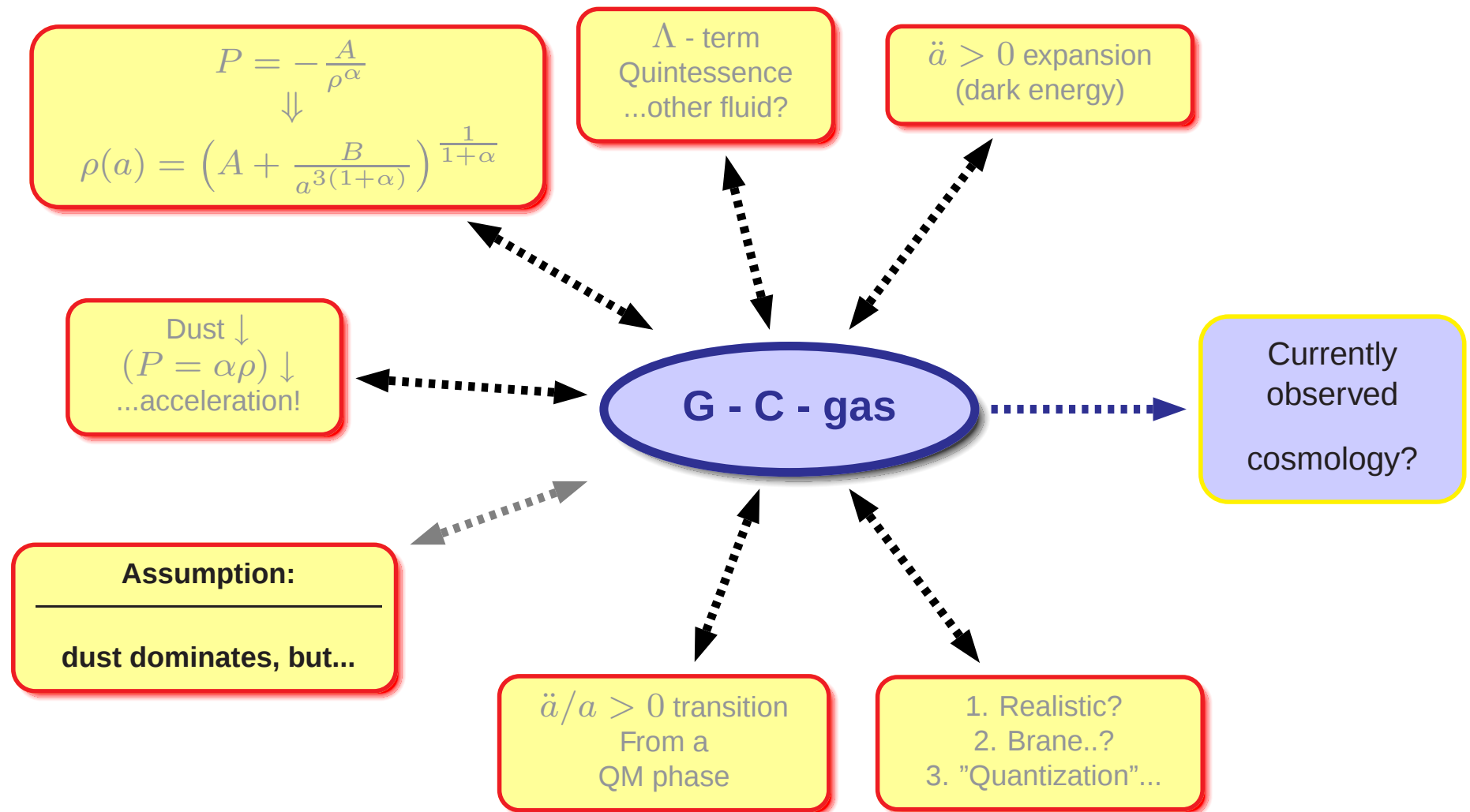
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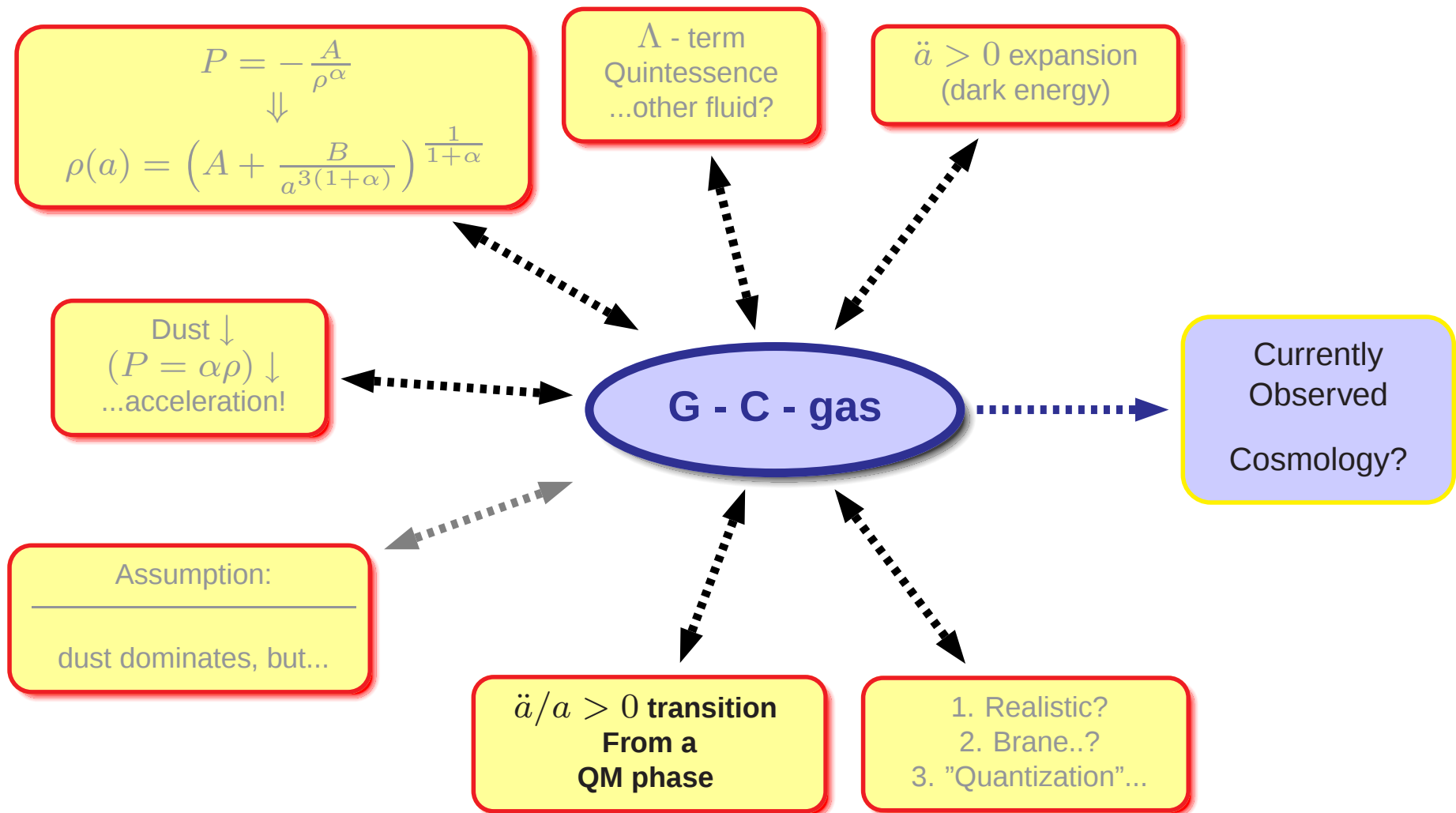
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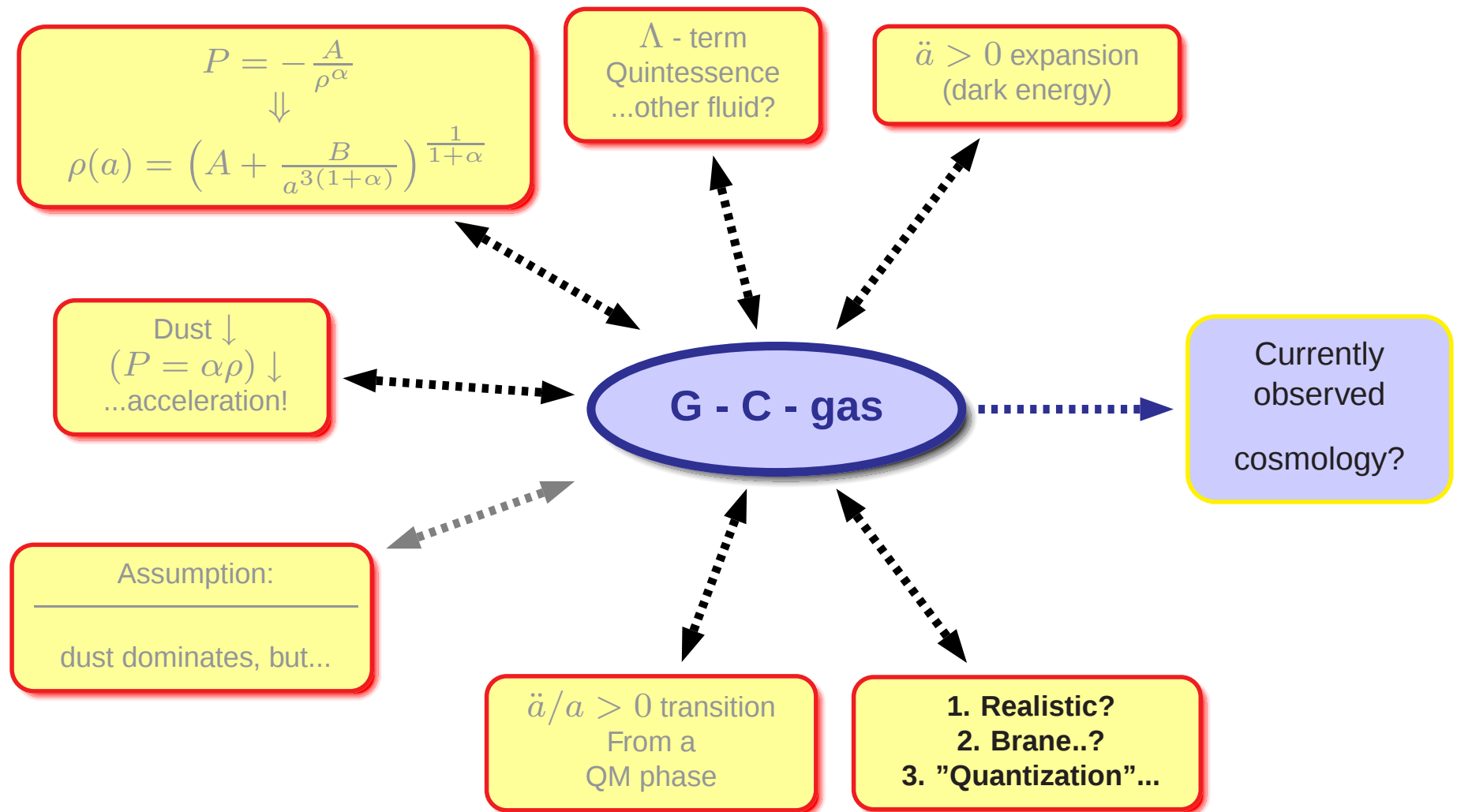
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Lagrangian:

$$L = -\frac{3\pi}{4G} \left(\frac{\dot{a}^2 a}{N} - Na \right) - 2\pi a^3 N \frac{\Lambda}{8\pi G} - 2\pi a^3 N \rho$$

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$$L = -\frac{3\pi}{4G} \left(\frac{\dot{a}^2 a}{N} - Na \right) - 2\pi a^3 N \frac{\Lambda}{8\pi G} - 2\pi a^3 N \rho$$

Wheeler-DeWitt equation:

$$\left[-\frac{G}{3\pi} \frac{d^2}{da^2} + \frac{3\pi}{4G} V_0(a) \right] \Psi(a) = 0$$

$$V_0(a) = a^2 - \lambda a^4 - \frac{8\pi G}{3} \rho(a) a^4$$

Semiclassical approximation $\Psi = c_1 \psi_1 + c_2 \psi_2$:

(in/out - going; in/de-creasing modes)

$$\psi_i = \exp \left[-\frac{1}{G} S_0(a) \right], \quad \left(\frac{dS_0(a)}{da} \right)^2 = \frac{9\pi^2}{4} V_0(a)$$

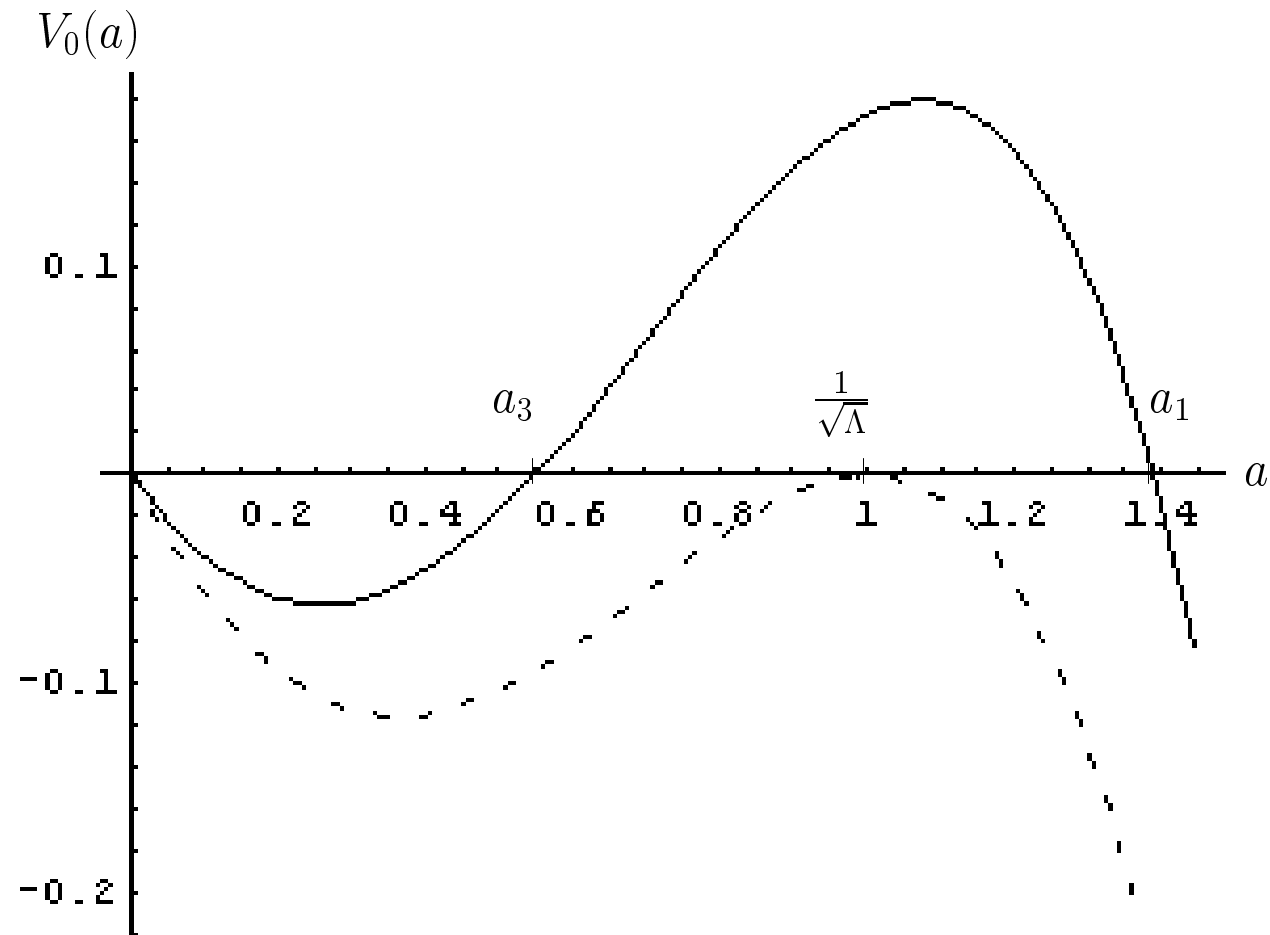
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Linear approximation $\rightarrow$ (saddle points):

$$V_0(a) \simeq a^2 - \lambda a^4 - a B^{\frac{1}{1+\alpha}}, \quad \Leftarrow a \ll \left[(\alpha + 1) \frac{B}{A} \right]^{\frac{1}{3(\alpha+1)}}$$

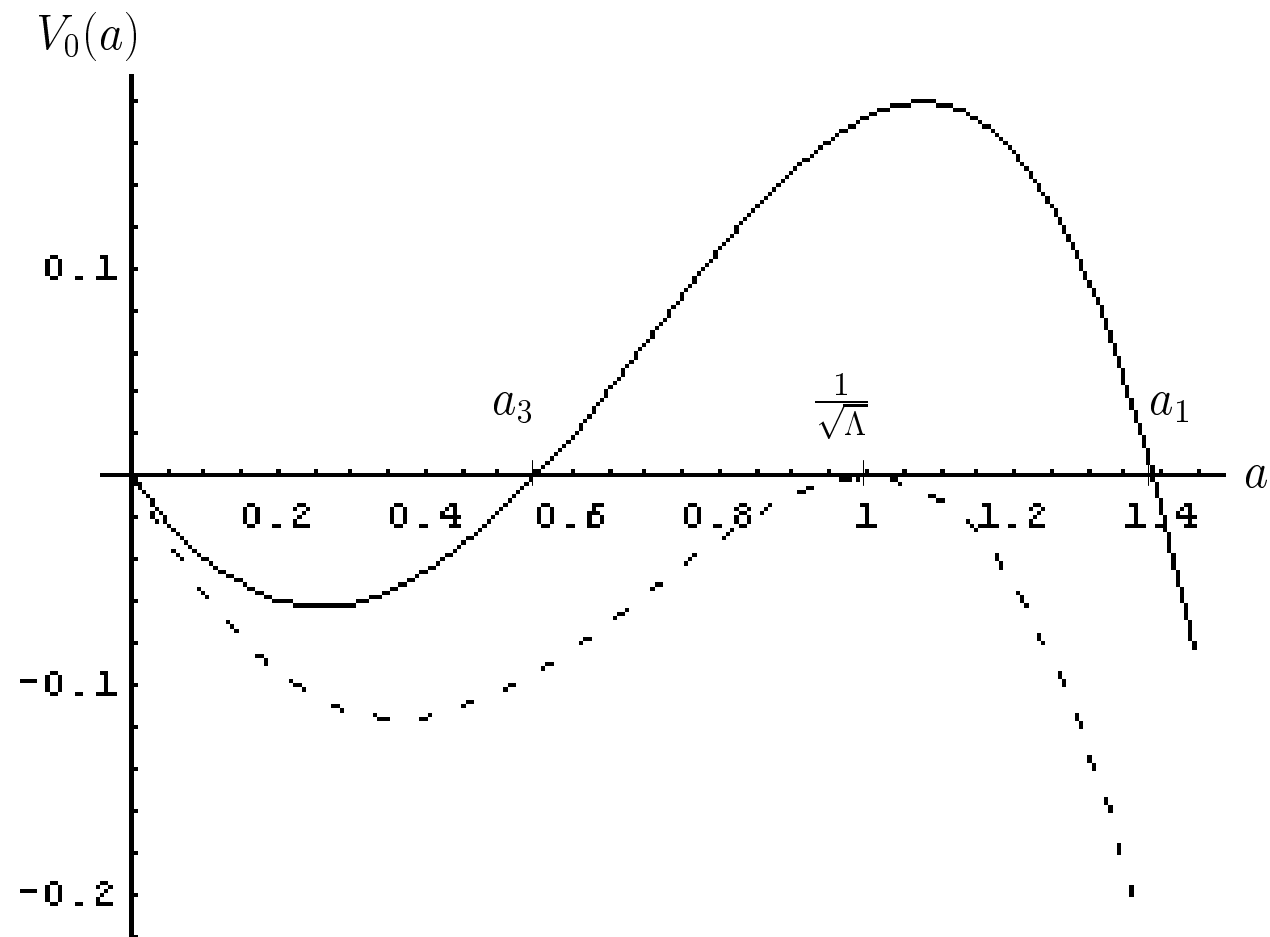


Turning point: depends on $\Lambda \tilde{B}^2$ where $\tilde{B} \equiv B^{\frac{1}{1+\alpha}}$

- $\Lambda \tilde{B}^2 > 4/9 \Rightarrow$ a classically allowed region.
- $\tilde{B} = \frac{2}{3} \frac{1}{\sqrt{\Lambda}} \Rightarrow$ two classically allowed regions.

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- $\tilde{B} = \frac{2}{3} \frac{1}{\sqrt{\Lambda}} \Rightarrow$ two classically allowed regions.
- ★ $\Lambda \tilde{B}^2 < 4/9 \Rightarrow$ two turning points a_3, a_1 ; two classically allowed regions ($0 < a < a_3, a_1 < a$) and one classically forbidden region ($a_3 < a < a_1$).



- FRW & $\phi = e^{i\theta} \varphi$; $m^2 \sim \varphi^2$; π_θ conserved $\rightarrow$ Born - Infeld phenomenology
- DeS 3+1 – brane in 4+1 Minkowski – bulk

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Go Quantum:

The wave function in the WKB approximation is related to $S_0(a)$ which depends on $\Lambda \tilde{B}^2$ and the range of values of a ; Compute:

$$\int \sqrt{|V_0(a)|} da \equiv \sqrt{\frac{\Lambda}{3}} \int I(a) da$$

$$\Lambda \tilde{B}^2 < 4/9$$

$$\begin{aligned} \int_{a_1}^a I(a) da &= \frac{aI}{3} - \beta \frac{a^2(a - a_1)}{I} + \gamma \left\{ a_1 a_3 (a_1^2 - a_3^2) \Pi \left(v, \frac{2a_1 + a_3}{2a_3 + a_1}, q \right) \right. \\ &\quad + a_1 a_3^2 (a_1 + a_3) F(v, q) - 6\beta^2 F(v, q) + \beta a_3^2 F(v, q) \\ &\quad \left. + \beta \frac{(a_1 + a_3)^2}{2a_3 + a_1} [(a_3 - a_1) \Pi(v, 1, q) + (2a_1 + a_3) F(v, q)] \right\} \end{aligned}$$

$$\begin{aligned} \text{where } \beta &\equiv \frac{(a_1^2 + a_1 a_3 + a_3^2)}{3}, \gamma \equiv \left(\sqrt{a_1(2a_3 + a_1)} \right)^{-1}, q \equiv \sqrt{\frac{a_3(2a_1 + a_3)}{a_1(2a_3 + a_1)}}, \\ v &\equiv \arcsin \sqrt{\frac{(2a_3 + a_1)(a - a_1)}{(2a_1 + a_3)(a - a_3)}}, \end{aligned}$$

$$\Lambda \tilde{B}^2 < 4/9$$

$$\begin{aligned} \int_a^{a_3} I(a) da = & -\frac{aI}{3} - \beta \frac{a^2(a_3 - a)}{I} + \gamma \left\{ a_1 a_3 (a_1 + a_3) \left[(a_3 - a_1) \Pi \left(\chi, \frac{a_3}{a_1}, q \right) \right. \right. \\ & - 6\beta^2 F(\chi, q) + \beta \gamma a_1^2 F(\chi, q) \\ & \left. \left. + \beta \frac{(a_1 + a_3)^2}{2a_3 + a_1} \left[(a_1 - a_3) \Pi(\chi, q^2, q) + (2a_3 + a_1) F(\chi, q) \right] \right] \right\} \end{aligned}$$

$$\text{with } \chi \equiv \arcsin \sqrt{\frac{a_1(a_3 - a)}{a_3(a_1 - a)}}$$

$$\Lambda \tilde{B}^2 < 4/9$$

$$\begin{aligned} \int_{a_3}^a I(a) da &= \frac{aI}{3} + \beta \frac{a^2(a - a_3)}{I} + \gamma \left\{ -a_1 a_3^2 (a_1 + a_3) \Pi \left(\delta, \frac{a_1 - a_3}{a_1}, r \right) \right. \\ &\quad + 6\beta^2 F(\delta, r) - \beta a_1 [(a_1 - a_3) F(\delta, r) + a_3 \Pi(\delta, 1, r)] \\ &\quad \left. + \beta (a_1 + a_3) [a_3 \Pi(\delta, r^2, r) - (2a_3 + a_1) F(\delta, r)] \right\} \end{aligned}$$

$$\text{with } \delta \equiv \arcsin \sqrt{\frac{a_1(a - a_3)}{(a_1 - a_3)a}} \text{ and } r \equiv \sqrt{\frac{a_1^2 - a_3^2}{(2a_3 + a_1)a_1}}$$

Hartle-Hawking wave function

$$\Psi(a) = C \cos \left[\frac{3\pi}{2G} \int_a^{a_3} \sqrt{-V_0(a)} da + \frac{\pi}{4} \right], \quad (a < a_3)$$

$$\Psi(a) = C \exp \left[\frac{3\pi}{2G} \int_{a_3}^a \sqrt{V_0(a)} da \right], \quad (a_3 < a < a_1)$$

$$\Psi(a) = 2\tilde{C} \cos \left[\frac{3\pi}{2G} \int_{a_1}^a \sqrt{-V_0(a)} da - \frac{\pi}{4} \right], \quad (a_1 < a)$$

where $\tilde{C} \equiv C \exp \left[\frac{3\pi}{2G} \int_{a_3}^{a_1} \sqrt{V_0(a)} da \right]$ and C is an arbitrary constant.

Vilenkin wave function

$$\begin{aligned}\Psi(a) &= C \exp \left[-\frac{3\pi}{2G} \int_{a_3}^{a_1} \sqrt{V_0(a)} da \right] \cos \left[\frac{3\pi}{2G} \int_a^{a_3} \sqrt{-V_0(a)} da + \frac{\pi}{4} \right] \\ &+ i4C \exp \left[\frac{3\pi}{2G} \int_{a_3}^{a_1} \sqrt{V_0(a)} da \right] \cos \left[\frac{3\pi}{2G} \int_a^{a_3} \sqrt{-V_0(a)} da - \frac{\pi}{4} \right], \\ &\hspace{15em} (a < a_3)\end{aligned}$$

$$\begin{aligned}\Psi(a) &= C \left(\exp \left[-\frac{3\pi}{2G} \int_a^{a_1} \sqrt{V_0(a)} da \right] + \right. \\ &\left. + 2i \exp \left[\frac{3\pi}{2G} \int_a^{a_1} \sqrt{V_0(a)} da \right] \right), \quad (a_3 < a < a_1)\end{aligned}$$

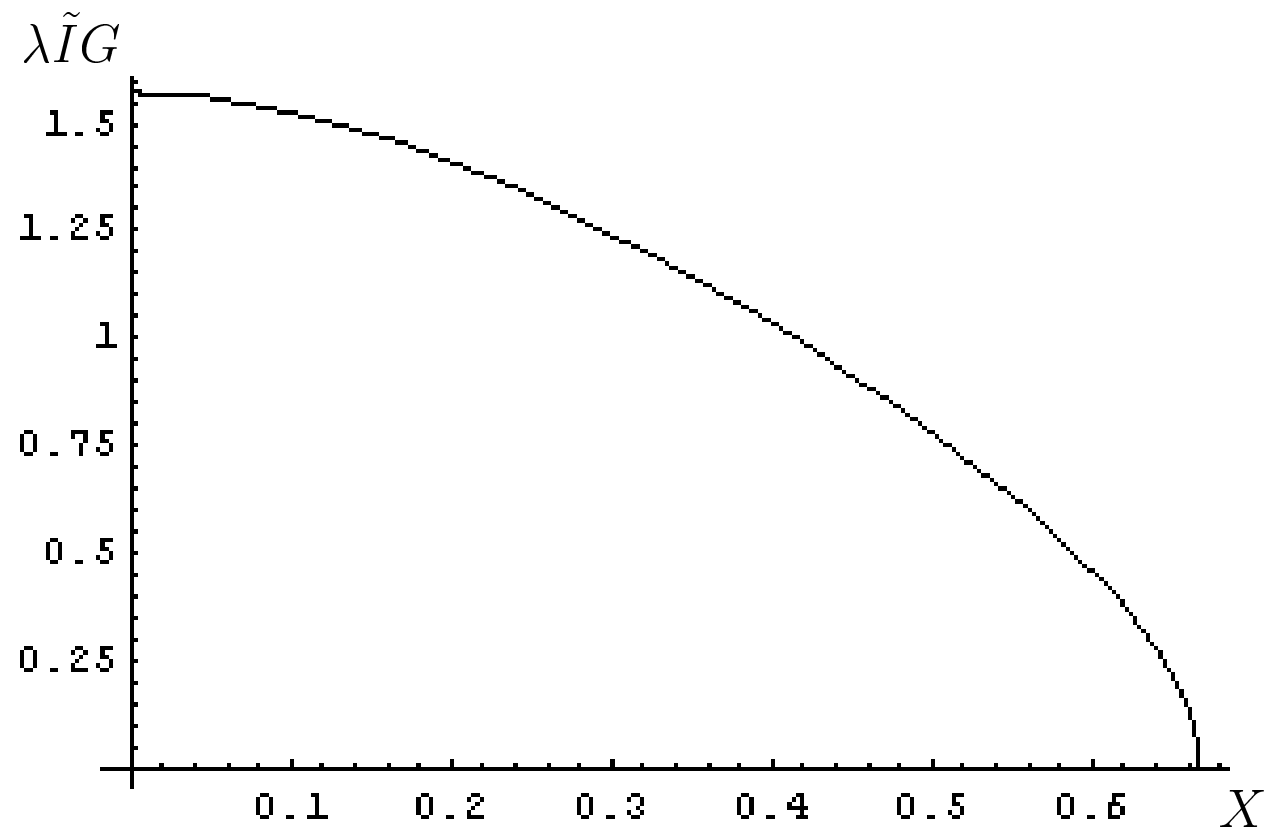
$$\Psi(a) = 2C e^{i\pi/4} \exp \left[-i \frac{3\pi}{2G} \int_{a_1}^a \sqrt{-V_0(a)} da \right], \quad (a_1 < a)$$

Transition amplitude

- Hartle-Hawking wave function: $A = \exp(2\tilde{I})$

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- Hartle-Hawking wave function: $A = \exp(2\tilde{I})$
- Vilenkin wave function: $A = -\exp(2\tilde{I})$



1. When $X \equiv B^{1/(\alpha+1)} \sqrt{\Lambda}$ increases $\lambda \tilde{I} G$ decreases (for fixed values of Λ).

1a. When X increases $\Rightarrow$ The transition amplitude for the Vilenkin boundary condition increases, while the transition amplitude for the Hartle-Hawking condition decreases (for fixed values of Λ).

1. When $X \equiv B^{1/(\alpha+1)}\sqrt{\Lambda}$ increases $\lambda\tilde{I}G$ decreases (for fixed values of Λ).
 - 1a. When X increases $\Rightarrow$ The transition amplitude for the Vilenkin boundary condition increases, while the transition amplitude for the Hartle-Hawking condition decreases (for fixed values of Λ).
2. When B increases X increases (for fixed values of Λ and α).
3. When α increases X can increase or decrease (for fixed values of Λ and B).

There is a **singularity** at $a = 0$ (R diverges)

$$R = 4 \left\{ \Lambda + 2\pi G \left[A + \frac{B}{a^{3(1+\alpha)}} \right]^{-\frac{\alpha}{\alpha+1}} \left[4A + \frac{B}{a^{3(1+\alpha)}} \right] \right\}.$$

There is a **singularity** at $a = 0$ (R diverges)

$$R = 4 \left\{ \Lambda + 2\pi G \left[A + \frac{B}{a^{3(1+\alpha)}} \right]^{-\frac{\alpha}{\alpha+1}} \left[4A + \frac{B}{a^{3(1+\alpha)}} \right] \right\}.$$

We **assume** that the wave function **vanishes** at the origin:

$$\Psi(a = 0) = 0$$

Parameters α and B restricted to a set of quantized values

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For the Hartle-Hawking wave function, e.g.,:

$$\frac{3\pi}{2G} \int_0^{a_3} \sqrt{-V_0(a)} da - \frac{\pi}{4} = n\pi, \quad n \in \mathbb{Z}$$

$$\Rightarrow \frac{3\pi}{4G} B^{2/(1+\alpha)} - 1 \simeq 4n$$

Next order of magnitude - going beyond WKB...

$$\left(A + \frac{B}{a^{3(1+\alpha)}} \right)^{\frac{1}{1+\alpha}} \approx \frac{B^{\frac{1}{1+\alpha}}}{a^3} \left[1 + \frac{1}{1+\alpha} \frac{A}{B} a^{3(\alpha+1)} + \frac{1}{2} \frac{1}{1+\alpha} \left(\frac{1}{1+\alpha} - 1 \right) \frac{A^2}{B^2} a^{6(\alpha+1)} + \dots \right].$$

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Wheeler-DeWitt equation

$$\left[-\frac{G}{3\pi} \frac{d^2}{da^2} + \frac{3\pi}{4G} V(a) \right] \Psi(a) = 0,$$

Wave function

$$\Psi(a) = \exp \left[-\frac{1}{G} \left(S_0(a) + A S_1(a) + A^2 S_2(a) + \dots \right) \right]$$

Hamilton-Jacobi eq and ...

$$\left(\frac{dS_0(a)}{da} \right)^2 = \frac{9\pi^2}{4} V_0(a),$$
$$\frac{2}{3\pi G} \frac{dS_0(a)}{da} \frac{dS_1(a)}{da} - \frac{1}{3\pi} \frac{d^2 S_1(a)}{da^2} + \frac{3\pi}{4G} V_1(a) = 0,$$

Hamilton-Jacobi eq and ...

$$\left(\frac{dS_0(a)}{da} \right)^2 = \frac{9\pi^2}{4} V_0(a),$$

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Potential terms

$$V_0(a) = a^2 - \lambda a^4 - a B^{\frac{1}{1+\alpha}},$$

$$V_1(a) = -\frac{a^{4+3\alpha}}{2+\alpha} B^{-\frac{\alpha}{1+\alpha}},$$

Compute wave function

$$J = \int_0^{a_3} \frac{V_1(a)}{\sqrt{-V_0(a)}} da,$$

$$J \approx B^{4+3\alpha} \mathbf{B} \left(9/2 + 3\alpha, \frac{1}{2} \right),$$

where $\mathbf{B}$ is the beta function ↓

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where $\mathbf{B}$ is the beta function $\downarrow$

$$\Psi(a) \approx \exp \left[-\frac{1}{G} (S_0(a) + AS_1(a)) \right]$$

