CLASSICAL MEASUREMENTS IN CURVED SPACE-TIMES

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THE TASK OF THE BOOK

 A physical measurement is meaningful only if one identifies in a non ambiguous way who is the observer and what is being observed. The same observable can be the target of more than one observer so we need a suitable algorithm to compare their measurements. This is the task of the theory of measurement which we develop here in the framework of general relativity.

THE MEASUREMENT PROTOCOL

- Identify the covariant equations which describe the phenomenon under investigation.
- Identify the observer who makes the measurements.
- (iii) Chose a frame adapted to that observer, allowing the space-time to be split into the observer's space and time.
- (iv) Decide whether the intended measurement is local or non-local with respect to the background curvature.
- (v) Identify the frame components of those quantities which are the observational targets.
- (vi) Find a physical interpretation of the above components, following a suitable criterion such as a comparison with what is known from special relativity or from non-relativistic theories.
- (vii) Verify the degree of residual ambiguity in the interpretation of the measurements and decide on a strategy to eliminate it.

$$\frac{D^2 Y^{\alpha}}{d\tau_U^2} = -R^{\alpha}{}_{\beta\gamma\delta} U^{\beta} Y^{\gamma} U^{\delta} + Y^{\sigma} \nabla_{\sigma} a(U)^{\alpha}.$$

WITH RESPECT TO A REFENCE FRAME

$$\ddot{Y}^{\hat{a}} + \mathcal{K}_{(U,E)}{}^{\hat{a}}{}_{\hat{b}}Y^{\hat{b}} = 0,$$

TIDAL MATRIX

$$\mathcal{K}_{(U,E)}{}^{\hat{a}}{}_{\hat{b}} = [T_{(\text{fw},U,E)} - S(U) + \mathcal{E}(U)]^{\hat{a}}{}_{\hat{b}}.$$

TWIST TENSOR

$$\begin{split} T_{(\text{fw},\,U,\,E)}{}^{\hat{a}}{}_{\hat{b}} &= \dot{C}_{(\text{fw})}{}^{\hat{a}}{}_{\hat{b}} - [C^2_{(\text{fw})}]^{\hat{a}}{}_{\hat{b}} - 2C_{(\text{fw})}{}^{\hat{a}}{}_{\hat{c}} \, k(U)^{\hat{c}}{}_{\hat{b}} \\ &= \delta^{\hat{a}}_{\hat{b}} \zeta^2_{(\text{fw})} - \zeta^{\hat{a}}_{(\text{fw})} \zeta_{(\text{fw})\hat{b}} - \epsilon^{\hat{a}}{}_{bf} \dot{\zeta}^{\hat{f}}_{(\text{fw})} \\ &- 2\epsilon^{\hat{a}}{}_{\hat{f}\hat{c}} \zeta^{\hat{f}}_{(\text{fw})} k(U)^{\hat{c}}{}_{\hat{b}}, \end{split}$$

FERMI-WALKER STRAIN TENSOR

$$S(U) = \nabla(U)a(U) + a(U) \otimes a(U),$$

GRAVITOELECTRIC TENSOR

$$R(U,Y)U \equiv R^{\alpha}{}_{\beta\gamma\delta}U^{\beta}Y^{\gamma}U^{\delta} = \mathcal{E}(U)^{\alpha}{}_{\gamma}Y^{\gamma}$$

THE RELATIVE ACCELERATION EQUATION IN TERMS OF TETRAD COMPONENTS

$$\frac{d^2Y^{\hat{r}}}{d\tau^2} = \left\{ -\mathcal{E}(U)_{\hat{r}\hat{r}} + \partial_{\hat{r}}a(U)_{\hat{r}} + (a(U)_{\hat{r}})^2 + (\zeta_{(\text{fw})_{\hat{\theta}}})^2 \right\} Y^{\hat{r}},$$

$$\frac{d^2Y^{\hat{\theta}}}{d\tau^2} = \left\{ -\mathcal{E}(U)_{\hat{\theta}\hat{\theta}} + \partial_{\hat{\theta}}a(U)_{\hat{\theta}} + \frac{1}{2}(E^{\theta}_{\hat{\theta}})^2(\partial_{\hat{r}}g_{\theta\theta})a(U)_{\hat{r}} \right\} Y^{\hat{\theta}},$$

$$\frac{d^2Y^{\hat{\phi}}}{d\tau^2} = \left\{ -\mathcal{E}(U)_{\hat{\phi}\hat{\phi}} - \Gamma_{\hat{r}\hat{\phi}\hat{\phi}}a(U)_{\hat{r}} + (\zeta_{(\text{fw})\hat{\theta}})^2 \right\} Y^{\hat{\phi}},$$

THE CHOICE OF THE METRIC

$$\begin{split} ds^2 &= -\left(1 - \frac{2\mathcal{M}r}{\Sigma}\right)dt^2 - \frac{4a\mathcal{M}r\sin^2\theta}{\Sigma}dtd\phi + \frac{A}{\Sigma}\sin^2\theta d\phi^2 \\ &+ \frac{\Sigma}{\Delta}dr^2 + \Sigma d\theta^2, \end{split}$$

THE CHOICE OF THE OBSERVER

$$U = \Gamma(\partial_t + \zeta \partial_\phi) ,$$

$$\Gamma = \left[1 - \frac{2\mathcal{M}r}{\Sigma} (1 - a\zeta \sin^2 \theta)^2 - (r^2 + a^2)\zeta^2 \sin^2 \theta \right]^{-1/2}.$$

THE CHOICE OF THE REFERENCE FRAME

$$E_{\hat{t}} = U = \Gamma(\partial_t + \zeta \partial_\phi),$$

$$E_{\hat{r}} = (\Delta/\Sigma)^{1/2} \partial_r,$$

$$E_{\hat{\theta}} = \Sigma^{-1/2} \partial_\theta,$$

$$E_{\hat{\phi}} = \bar{\Gamma}(\partial_t + \bar{\zeta} \partial_\phi),$$

THE STRESS TENSOR

$$a(U)^{\hat{r}} = \frac{\Gamma^2 \sqrt{\Delta}}{\sqrt{\Sigma}} \left[\frac{\mathcal{M}(r^2 - a^2 \cos^2 \theta)}{\Sigma^2} (1 - a\zeta \sin^2 \theta)^2 - r\zeta^2 \sin^2 \theta \right],$$

$$a(U)^{\hat{\theta}} = -\frac{\Gamma^2 \sin \theta \cos \theta}{\sqrt{\Sigma}} \left[\frac{2\mathcal{M}r}{\Sigma^2} [(r^2 + a^2)\zeta - a]^2 + \Delta \zeta^2 \right].$$

THE OUTCOME OF A MEASUREMENT

THE DISCOVERY OF A NEW EFFECT:

THE RELATIVISTIC THRUST ANOMALY

A NEW GENERAL RELATIVISTIC EFFECT

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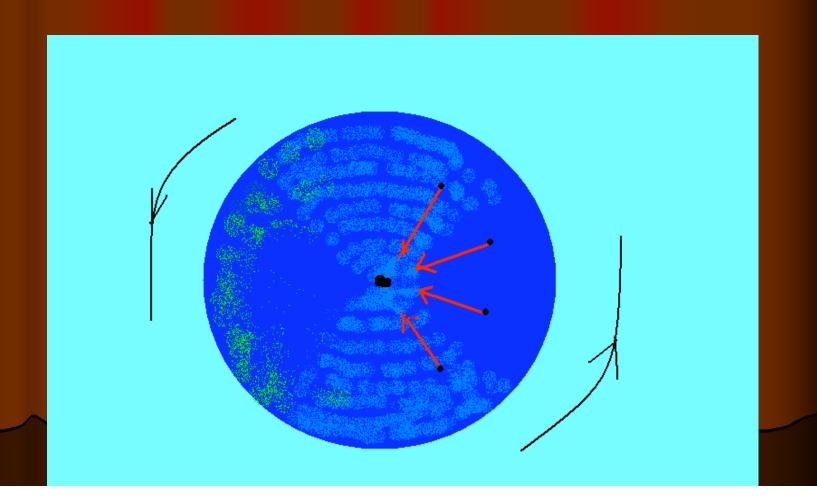
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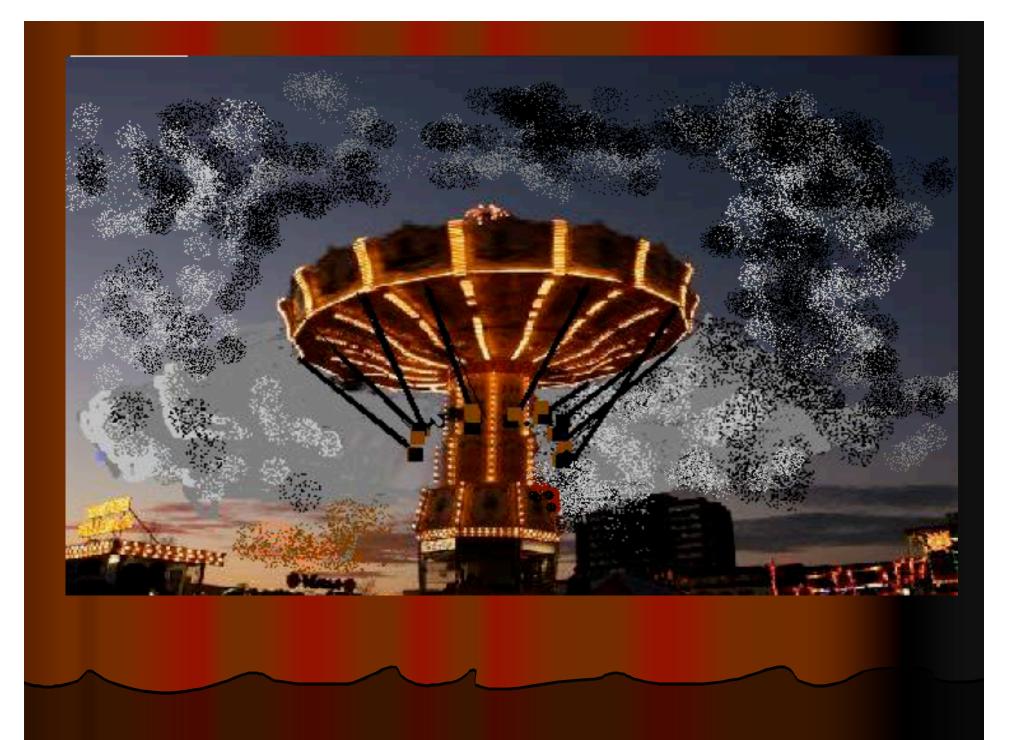
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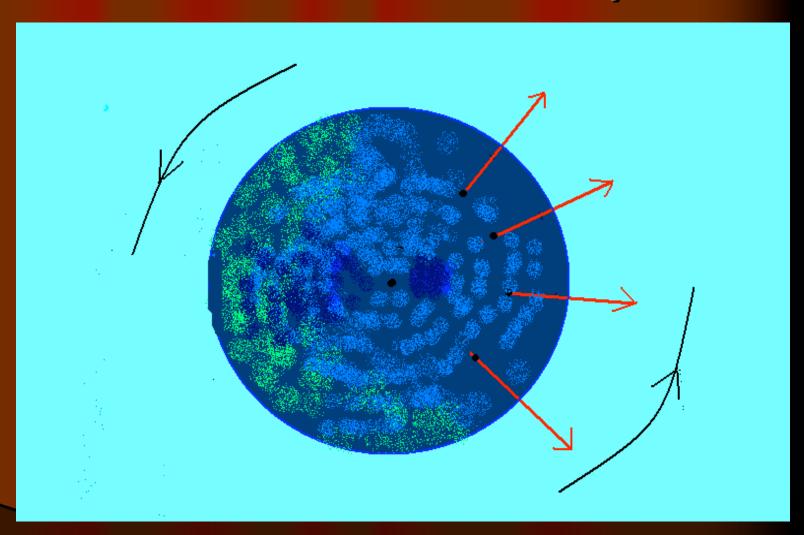
CENTRIFUGAL FORCE

Normal behaviour: the thrust acts inwardly





The thrust acts outwardly The thrust anomaly



IT MANIFESTS IN TWO SITUATIONS

In Schwarzschild at 2m < r < 3 M

Abramowicz

Acta Phys. Pol. B5 (1974) 327

In Kerr at $r_+ < r < 00$

de Felice

1994

Circular orbits in the space-time of a rotating source

$$y = \frac{\zeta}{1 - a\zeta}$$

$$\Delta = r^2 + a^2 - 2\mathcal{M}r$$

$$y_{K\pm} \equiv \pm \sqrt{\frac{\mathcal{M}}{r^3}}$$

$$y_{c_{\pm}} = (a \pm \sqrt{\Delta})/r^2$$

Specific thrust acting on a general nongeodesic equatorial spatially circular orbit in the space-time of a rotating source

$$a(U) = \frac{\Delta^{1/2}}{r^2} \frac{(y - y_{k_+})(y - y_{k_-})}{(y - y_{c_+})(y - y_{c_-})},$$

$$\theta = \pi/2$$

Gradient of the specific thrust with respect to the angular velocity y

$$\frac{\partial a(U)}{\partial y}\bigg|_{r} = -\frac{2a\Delta^{1/2}}{r^4} \frac{(y - y_{0_{+}})(y - y_{0_{-}})}{(y - y_{c_{+}})^2(y - y_{c_{-}})^2}$$

THE MEASUREMENT OF THE THRUST

$$a(U_{(\mathrm{crit})-}) = \frac{\Delta^{1/2}}{r^2} \frac{\sqrt{1 - \frac{4\mathcal{M}a^2}{r^3} \left(1 - \frac{3\mathcal{M}}{r}\right)^{-2}} - 1}{\sqrt{1 - \frac{4\mathcal{M}a^2}{r^3} \left(1 - \frac{3\mathcal{M}}{r}\right)^{-2}} - \frac{2a^2}{r^2} \left(1 - \frac{3\mathcal{M}}{r}\right)^{-1} - 1}}.$$

Values of y where the specific thrust has a maximum

$$y_{0\pm} = -\frac{1}{2a} \left[1 - \frac{3\mathcal{M}}{r} \mp \sqrt{\left(1 - \frac{3\mathcal{M}}{r}\right)^2 - \frac{4\mathcal{M}a^2}{r^3}} \right]$$

Behaviour of the specific thrust as function of *y* at a fixed radius

